



Original Article

The character table of a subgroup $2^7:G_2(2)$ of $Sp_8(2)$

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ABSTRACT: In this paper, the ordinary character table of a finite extension of structure $\bar{G} = 2^7:G_2(2)$ is computed via the Fischer-Clifford matrices technique. The group \bar{G} sits maximally in the affine subgroup $2^7:Sp_6(2)$ of the symplectic group $Sp_8(2)$.

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1. Introduction

The Chevalley group $G_2(2)$ of type 2 has an absolutely irreducible representation of dimension 6 over $GF(2)$ and hence the split extension $\bar{S} = 2^6:G_2(2)$ exists [22]. In fact, \bar{S} sits maximally in the maximal subgroup $\bar{S}_1 = 2^6:Sp(6, 2)$ of the sporadic simple group Fi_{22} [5]. Moreover, the Schur Multiplier $M(\bar{S})$ of \bar{S} is isomorphic to the cyclic group \mathbb{Z}_2 of order 2 and we obtain the 2-cover group $2.\bar{S}$ of \bar{S} which is isomorphic to a group of structure $\bar{G} = 2^7:G_2(2)$. Furthermore, \bar{G} also sits maximally in the 2-cover $2.\bar{S}_1$ of \bar{S}_1 which in turn is a maximal subgroup of $2.Fi_{22}$. Note that $2.\bar{S}_1 \cong ASp_8(2) = 2^7:Sp_6(2)$, where $ASp_8(2)$ is the affine subgroup of the symplectic group $Sp_8(2)$. Therefore, $Sp_8(2)$ contains an isomorphic copy of \bar{G} . The information above is verified with GAP [20] and we state it as Theorem 1.1.

Theorem 1.1. *The cover groups \bar{G} and $2.\bar{S}_1$ of \bar{S} and \bar{S}_1 , respectively, are subgroups of $Sp_8(2)$.*

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In the online ATLAS [22], it can be seen that $Sp_8(2)$ has an absolutely irreducible representation of dimension 16 over $GF(2)$ and hence the split extension $\overline{G}_1 = 2^{16}:Sp_8(2)$ can be constructed. Using the Fischer-Clifford matrices technique [6], the ordinary irreducible character table of \overline{G}_1 was constructed in [14]. The group \overline{G} is one of the inertia factors of the action of \overline{G}_1 on the irreducible characters $\text{Irr}(2^{16})$ of 2^{16} . Therefore, the set $\text{Irr}(\overline{G})$ plays an important role in the construction of the ordinary character table of \overline{G}_1 and for this reason we consider \overline{G} as a subgroup of $Sp_8(2)$ and compute its ordinary character table using the Fischer-Clifford matrices technique. A Fischer-Clifford matrix (defined in Subsection 3.1) of \overline{G} , denoted by $M(g)$, is a square matrix with complex entries and is constructed for each conjugacy class $[g]$ of $Sp_8(2)$. In practice, we do not use the formal definition of a Fischer-Clifford matrix $M(g)$ to compute the character table of a finite extension group, for example \overline{G} , but instead use the powerful arithmetical properties (see Subsection 3.1) associated with $M(g)$. These properties are naturally inherited from the row and column orthogonal relations of the ordinary character table of a finite extension group. Therefore, we can compute the ordinary character tables of complicated finite extensions of p -groups with great efficiency using the Fischer-Clifford matrices technique. Since \overline{G} is a split extension with its kernel 2^7 a 7-dimensional $G_2(2)$ -module, the Fischer-Clifford matrices technique is an appropriate choice to compute the ordinary character table of \overline{G} .

From now on, we let \overline{G} be the split extension of the elementary abelian 2-group $P = 2^7$ by $G = G_2(2)$. The conjugacy classes of \overline{G} will be computed from each coset $Pg \in \frac{\overline{G}}{P}$, where g is a conjugacy class $[g]_G$ representative of G . The method that will be used for this purpose is called the coset analysis technique (see [10, 11]). MAGMA routines in [1], based on the coset analysis technique, are used to compute the classes of \overline{G} and their p power maps. The advantage of this routines is that we only need the generators of G as matrices of degree 7 over $GF(2)$ and identify P with a 7-dimensional vector space $V_7(2)$ over $GF(2)$. We do not require a representation of the whole of \overline{G} . Having the classes of \overline{G} in coset analysis format, we can proceed to construct the Fischer-Clifford matrices $M(g)$ of \overline{G} corresponding to each class representative $g \in G$. The ordinary character table of \overline{G} is then constructed using the matrices $M(g)$ and the ordinary character tables of the so-called inertia factors H_i which are subgroups of G . The character table of \overline{G} will be partitioned into blocks according to each inertia group \overline{H}_i of $\text{Irr}(P)$ in \overline{G} . See Subsection 3.1 for the definitions of \overline{H}_i and H_i .

The character table of a nonsplit extension $\overline{G}_2 = 2^7:G_2(2)$ was computed in [18] by the Fischer-Clifford matrices technique. Although, the ordinary character tables of \overline{G} and \overline{G}_2 do not coincide, but they have the same number of irreducible characters. Computations are done in GAP and MAGMA [4]. ATLAS notation [5] is used, unless stated otherwise.

2. Conjugacy Classes of \overline{G}

2.1. Construction of G as a 7×7 Matrix Group over $GF(2)$

A finite group F is called a $n \times n$ matrix group over a finite field K if each element of F can be represented as an invertible $n \times n$ matrix with entries in K . In [1], the symplectic group $Sp_6(2) \leq ASp_8(2)$ was constructed as a 7×7 matrix group over $GF(2)$. Since $G \leq Sp_6(2)$, we construct G as a 7×7 matrix group over $GF(2)$ within $Sp_6(2)$ and its generators g_1 and g_2 of orders of 2 and 6 are listed in Figure 1.

$$g_1 = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 1 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 1 & 0 & 1 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix} \quad g_2 = \begin{pmatrix} 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 1 & 1 & 0 & 1 & 0 & 0 & 0 \\ 1 & 1 & 1 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 1 & 0 \\ 0 & 1 & 1 & 1 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 & 1 & 1 & 1 \\ 1 & 0 & 0 & 1 & 1 & 1 & 1 \end{pmatrix}$$

Figure 1: Generators of G

2.2. Actions of G on P and $\text{Irr}(P)$

The linear group $G = \langle g_1, g_2 \rangle$ acts on P , where we regard P as a 7-dimension vector space $V_7(2)$ over $GF(2)$, and obtain 4 orbits of lengths 1, 1, 63 and 63 with corresponding stabilizers $P_1 = G$, $P_2 = G$, $P_3 = 4^2:D_{12}$ and $P_4 = 4^2:D_{12}$. The permutation character $\chi(G|P) = 1aaaa + 14aa + 21aa + 27aa$ of G on the classes of P is computed in terms of the ordinary irreducible characters of G (see character table of $G_2(2)$ either in ATLAS or in GAP library). Evaluating $\chi(G|P)(g)$ on a class representative $g \in G$ gives the number of elements of P fixed by an element g . Since G has 4 orbits on P it follows by Brauer's Theorem [7] that G also has 4 orbits on $\text{Irr}(P)$. To act G on the dual space P^* of P can be identified with the action of G on $\text{Irr}(P)$. This action resulted in four orbits

of lengths 1, 28, 36 and 63 with corresponding stabilizers $H_1 = G$, $H_2 = 3_+^{1+2}:8:2$, $H_3 = L_2(7):2$ and $H_4 = 4^2:D_{12}$ which are the inertia factors of \overline{G} on $\text{Irr}(P)$. Table 1 contains a summary of the information pertaining to the actions of G on P and $\text{Irr}(P)$. The fusion maps of the classes of the inertia factors H_i into the classes of G are found in Table 2. For example, the fusion map of H_4 into G indicates that the elements of order 2 contain in the conjugacy class $2D$ of H_4 belong to the elements of order 2 in the conjugacy class $2B$ of G .

Table 1: Actions of G on P and $\text{Irr}(P)$

	Action of G on P	Action of G on $\text{Irr}(P)$
Number and lengths of Orbits	$ O_1 = 1$	$ O_1 = 1$
	$ O_2 = 1$	$ O_2 = 28$
	$ O_3 = 63$	$ O_3 = 36$
	$ O_4 = 63$	$ O_4 = 63$
Structure of point stabilizers	$P_1 = G_2(2)$	$H_1 = G_2(2)$
	$P_2 = G_2(2)$	$H_2 = 3_+^{1+2}:8:2$
	$P_3 = 4^2:D_{12}$	$H_3 = L_2(7):2$
	$P_4 = 4^2:D_{12}$	$H_4 = 4^2:D_{12}$
Size of stabilizers	$ P_1 = 12096$	$ H_1 = 12096$
	$ P_2 = 12096$	$ H_2 = 432$
	$ P_3 = 192$	$ H_3 = 336$
	$ P_4 = 192$	$ H_4 = 192$
Number of conjugacy classes $[g]$ of P_i and H_i	$ [g]_{P_1} = 16$	$ [g]_{H_1} = 16$
	$ [g]_{P_2} = 16$	$ [g]_{H_2} = 14$
	$ [g]_{P_3} = 14$	$ [g]_{H_3} = 9$
	$ [g]_{P_4} = 14$	$ [g]_{H_4} = 14$

Table 2: The fusion maps of H_i into G

$ C_{H_2}(h) $	$[h]_{H_2} \rightarrow [g]_G$	$ C_{H_2}(h) $	$[h]_{H_2} \rightarrow [g]_G$
432	1A 1A	24	6A 6A
48	2A 2A	6	6B 6B
12	2B 2B	8	8A 8A
216	3A 3A	8	8B 8A
18	3B 3B	12	12A 12C
24	4A 4A	12	12B 12A
12	4B 4B	12	12C 12B
$ C_{H_3}(h) $	$[h]_{H_3} \rightarrow [g]_G$	$ C_{H_3}(h) $	$[h]_{H_3} \rightarrow [g]_G$
336	1A 1A	6	6A 6B
16	2A 2A	7	7A 7A
12	2B 2B	8	8A 8B
6	3A 3B	8	8B 8B
8	4A 4A	12	
$ C_{H_4}(h) $	$[h]_{H_4} \rightarrow [g]_G$	$ C_{H_4}(h) $	$[h]_{H_4} \rightarrow [g]_G$
192	1A 1A	32	4A 4A
64	2A 2A	32	4B 4C
48	2B 2B	16	4C 4B
16	2C 2A	16	4D 4C
16	2D 2B	6	6A 6B
16	2E 2B	8	8A 8A
6	3A 3B	8	8B 8B

2.3. Coset Analysis Technique

The conjugacy classes of \overline{G} is computed by the method of coset analysis (see [10, 11, 12]). In this subsection, we give a brief description of the coset analysis method for a split extension $SE = EA:Q$ of an elementary abelian p -group EA of order p^n by a linear matrix group Q of degree n over the field $GF(p)$. The group EA is regarded as a vector space $V_n(p)$ of dimension n over the finite field $GF(p)$ and is a Q -module over $GF(2)$, where upon the matrix

group Q acts naturally. A coset $(EA)q$ is considered for each conjugacy class $[q]$ representative q in Q and then we consider the action of the stabilizer $C_q = EA:C_Q(q) = \{x \in SE | x((EA)q)x^{-1} = (EA)q\}$ of the coset EAq in SE by conjugation on the elements of $(EA)q$. Since C_q is split extension we will first act EA on $(EA)q$ to form k orbits Q_1, Q_2, \dots, Q_k , with each orbit Q_i containing $|EA|/k$ elements. Under the action of the centralizer $C_Q(q)$ of $q \in Q$, f_j of the k orbits Q_i fuse together to form an orbit O_j . The orbit O_j contains the elements from the coset $(EA)q$ which belong to a conjugacy class $[x_j]$ of SE with class representative x_j . Note that $\sum f_j = k$. The order of the centralizer $C_{SE}(x_j)$ of the class representative x_j is computed by $|C_{SE}(x_j)| = \frac{k|C_Q(q)|}{f_j}$. In this manner, from a coset $(EA)q$, we obtain a set of conjugacy classes $\bigcup_{j=1}^{c(q)} [x_j]$ of SE , with class representatives $X(q) = \{x_1, x_2, \dots, x_{c(q)}\}$.

We use a MAGMA routine, labelled as Programme A in [1], to compute the parameter f_j for our group \bar{G} . The value k is obtained by evaluate the permutation character $\chi(G|P)$, computed in Section 2.2, on a class representative $g \in G$. The MAGMA routine, Programme B in [1], is used to compute the orders of the elements of \bar{G} as well as their p -power maps (see column 9 of Table 3). All the information pertaining to the elements of \bar{G} is found in Table 3.

3. The Fischer-Clifford Matrices of \bar{G}

In this section, the Fischer-Clifford matrices technique will be applied to \bar{G} . A more general and detailed treatment of the Fischer-Clifford matrices technique is found in [1, 3, 6, 9, 12, 21]. For recent publications on the applications of the Fischer-Clifford matrices technique, see for example, [2, 13, 15, 16, 17].

3.1. General Construction of a Fischer-Clifford Matrix $M(g)$ of \bar{G}

From Section 2.2, \bar{G} has four orbits O_i on $\text{Irr}(P)$ with corresponding inertia groups $\bar{H}_i = P:H_i = \{x \in \bar{G} | \theta_i^x = \theta_i\}$, $i = 1, 2, 3, 4$, where $\theta_i \in O_i$ are representatives of the orbits O_i and $H_i \cong \frac{\bar{H}_i}{P}$ the inertia factor groups (see Table 1 above). Since P is elementary abelian, each θ_i extends to a $\psi_i \in \text{Irr}(\bar{H}_i)$, i.e. $\psi_i \downarrow_P = \theta_i$, by Mackey's Theorem (see Theorem 5.1.15 in [12]). Furthermore, by Theorem 5.1.7, Remark 5.1.8 and Theorem 5.1.19 in [12], an ordinary irreducible character $\chi = (\psi_i \bar{\beta})^{\bar{G}}$ of \bar{G} is obtained by induction of $\psi_i \bar{\beta} \in \text{Irr}(\bar{H}_i)$ to \bar{G} , where $P \subseteq \ker(\bar{\beta})$ of $\bar{\beta} \in \text{Irr}(\bar{H}_i)$. Note that $\bar{\beta}$ is a lifting of $\beta \in \text{Irr}(H_i)$ to \bar{H}_i . Hence, by Gallagher [8] we have Theorem 3.1 below for our group \bar{G} .

Theorem 3.1. $\text{Irr}(\bar{G}) = \bigcup_{i=1}^4 \{(\psi_i \bar{\beta})^{\bar{G}} | \bar{\beta} \in \text{Irr}(\bar{H}_i), P \subseteq \ker(\bar{\beta})\} = \bigcup_{i=1}^4 \{(\psi_i \bar{\beta})^{\bar{G}} | \beta \in \text{Irr}(H_i)\}$.

Therefore, the set $\text{Irr}(\bar{G})$ is partitioned into 4 blocks B_i with each block B_i corresponding to an inertia group \bar{H}_i and $|\text{Irr}(\bar{G})| = |\text{Irr}(H_1)| + |\text{Irr}(H_3)| + |\text{Irr}(H_3)| = 16 + 14 + 9 + 14 = 53$.

We define the set

$$R(g) = \{(i, y_k) | 1 \leq i \leq 4, H_i \cap [g] \neq \emptyset, 1 \leq k \leq r\},$$

where $y_k, k = 1, 2, \dots, r$, are representatives of conjugacy classes $[y_k]$ of H_i that fuse into a class $[g]$ of $H_1 = G$. Also, from the coset Pg , we have the set $X(g) = \{x_1, x_2, \dots, x_{c(g)}\}$ of class representatives of \bar{G} . Furthermore, we have the set of pre-images $\eta^{-1}(y_k) = \bigcup_{s=1}^{c(k)} \{y_{k_s}\}$ of a class representative $y_k \in H_i$ under the natural homomorphism $\eta: \bar{H}_i \rightarrow H_i$ such that y_{k_s} is conjugate to $x_j \in X(g)$. Then we evaluate $(\psi_i \bar{\beta})^{\bar{G}} \in \text{Irr}(\bar{G})$ on $x_j \in X(g)$ as in Theorem 3.2 below.

Theorem 3.2. $(\psi_i \bar{\beta})^{\bar{G}}(x_j) = \sum_{y_k: (i, y_k) \in R(g)} \left[\sum_{s=1}^{c(k)} \frac{|C_{\bar{G}}(x_j)|}{|C_{H_i}(y_{k_s})|} \psi_i(y_{k_s}) \right] \beta(y_k)$.

Proof. See [21] □

For a class representative $g \in G$, a Fischer-Clifford matrix $M(g) = (a_{(i, y_k)}^j)$ of \bar{G} is defined in Equation (1) below,

$$(a_{(i, y_k)}^j) = \left(\sum_{s=1}^{c(k)} \frac{|C_{\bar{G}}(x_j)|}{|C_{H_i}(y_{k_s})|} \psi_i(y_{k_s}) \right). \tag{1}$$

Whereas, Figure 2 represents $M(g)$ in matrix form.

Table 3: The Conjugacy Classes of \overline{G}

$[x]_G$	k	f_j	d	w	$[x]_{\overline{G}}$	$C_G(g)$	$ C_{\overline{G}}(x) $	2P	3P	5P	7P	$\rightarrow Sp_8(2)$
1A	128	1	(0,0,0,0,0,0)	(0,0,0,0,0,0)	1A	12096	1548288	1A	1A	1A	1A	1A
		1	(1,1,1,1,0,0,1)	(1,1,1,1,0,0,1)	2A		1548288	1A	2A	2A	2A	2A
		63	(0,1,0,1,0,1,1)	(0,0,0,0,1,0,0)	2B		24576	1A	2B	2B	2B	2B
		63	(1,0,0,0,0,0,0)	(1,0,0,0,0,0,0)	2C		24576	1A	2C	2C	2C	2C
2A	32	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	2D	192	6144	1A	2D	2D	2D	2H
		1	(1,1,1,1,0,0,1)	(0,0,0,0,0,0,0)	2E		6144	1A	2E	2E	2E	2I
		3	(0,1,0,0,0,0,0)	(0,0,0,0,0,0,0)	2F		2048	1A	2F	2F	2F	2G
		3	(1,0,0,0,0,0,1)	(0,0,0,0,0,0,0)	2G		2048	1A	2G	2G	2G	2J
		24	(0,0,1,1,1,0,1)	(1,0,0,0,0,1,0)	4A		256	2B	4A	4A	4A	4C
2B	16	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	2H	48	768	1A	2H	2H	2H	2N
		1	(1,1,1,1,0,1,0)	(0,0,0,0,0,0,0)	2I		768	2C	4B	4B	4B	4I
		1	(0,0,0,1,1,0,1)	(0,1,1,1,0,0,1)	4B		768	1A	2I	2I	2I	2O
		1	(1,1,1,1,0,0,1)	(1,0,1,0,0,0,0)	4C		768	2C	4C	4C	4C	4H
		6	(1,0,0,0,0,0,0)	(1,0,1,0,0,0,0)	4D		128	2C	4D	4D	4D	4J
		6	(1,1,0,0,1,1,0)	(0,1,1,1,0,0,1)	4E		128	2B	4E	4E	4E	4G
3A	2	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	3A	216	432	3A	1A	3A	3A	3B
		1	(1,1,1,1,0,0,1)	(1,1,1,1,0,0,1)	6A		432	3A	2A	6A	6A	6D
3B	8	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	3B	18	144	3B	1A	3B	3B	3C
		1	(1,1,1,1,0,0,1)	(0,1,1,1,0,0,1)	6B		144	3B	2A	6B	6B	6G
		3	(1,1,0,1,1,1,0)	(1,0,1,0,1,1,0)	6C		48	3B	2C	6C	6C	6E
		3	(1,0,0,0,0,0,0)	(0,1,1,1,1,0,0)	6D		48	3B	2B	6D	6D	6F
4A	8	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	4F	96	768	2D	4F	4F	4F	4L
		1	(0,1,0,0,0,1,0)	(0,0,0,0,0,0,0)	4G		768	2D	4G	4G	4G	4M
		6	(1,1,0,0,0,0,0)	(0,0,0,0,0,0,0)	4H		128	2F	4H	4H	4H	4K
4B	8	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	4I	48	384	2D	4I	4I	4I	4L
		1	(0,0,0,0,0,1,0)	(0,0,0,0,0,0,0)	4J		384	2D	4J	4J	4J	4M
		6	(0,0,0,0,0,0,1)	(0,0,0,0,0,0,0)	4K		64	2F	4K	4K	4K	4K
4C	8	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	4L	32	256	2D	4L	4L	4L	4X
		1	(0,1,0,0,0,1,0)	(0,0,0,0,0,0,0)	4M		256	2D	4M	4M	4M	4W
		2	(1,0,1,0,0,1,1)	(0,0,0,0,0,0,0)	4N		128	2F	4N	4N	4N	4V
		4	(1,1,1,0,0,0,1)	(0,0,0,0,0,0,0)	4O		64	2G	4O	4O	4O	4Y
6A	2	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	6E	24	48	3A	2D	6E	6E	6M
		1	(0,1,1,1,0,0,1)	(0,0,0,0,0,0,0)	6F		48	3A	2E	6F	6F	6N
6B	4	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	6G	6	24	3B	2H	6G	6G	6X
		1	(1,0,0,0,0,0,0)	(0,0,1,1,0,0,0)	12A		24	3B	2I	6H	6H	6W
		1	(1,1,1,0,0,0,1)	(0,0,0,0,0,0,0)	6H		24	6C	4B	12A	12A	12I
		1	(1,0,1,1,0,1,1)	(0,0,1,1,0,0,0)	12B		24	6C	4C	12B	12B	12J
7A	2	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	7A	7	14	7A	7A	7A	1A	7A
		1	(1,1,1,1,0,0,1)	(1,1,1,1,0,0,1)	14A		14	7A	14A	14A	2A	14A
8A	4	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	8A	8	32	4L	8A	8A	8A	8I
		1	(0,1,0,0,0,1,0)	(0,0,0,0,0,0,0)	8B		32	4L	8B	8B	8B	8J
		1	(1,1,1,0,0,0,1)	(0,0,0,0,0,0,0)	8C		32	4N	8D	8D	8C	8H
		1	(1,0,1,1,0,1,1)	(0,0,0,0,0,0,0)	8D		32	4N	8C	8C	8D	8H
8B	4	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	8E	8	32	4F	8E	8E	8E	8F
		1	(1,1,0,0,0,0,1)	(0,0,0,0,0,0,0)	8F		32	4F	8F	8F	8F	8G
		1	(0,0,1,1,0,0,0)	(0,0,0,0,0,0,0)	8G		32	4H	8G	8H	8H	8E
		1	(0,0,1,1,0,0,0)	(0,0,0,0,0,0,0)	8H		32	4H	8H	8G	8G	8E
12A	2	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	12C	12	24	6E	4F	12C	12C	12O
		1	(1,1,1,1,0,0,0)	(0,0,0,0,0,0,0)	12D		24	6E	4G	12D	12D	12P
12B	2	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	12E	12	24	6E	4I	12G	12E	12O
		1	(1,1,1,1,0,0,0)	(0,0,0,0,0,0,0)	12F		24	6E	4J	12H	12F	12P
12C	2	1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	12G	12	24	6E	4I	12E	12G	12O
		1	(0,0,0,0,0,0,0)	(0,0,0,0,0,0,0)	12H		24	6E	4J	12F	12H	12P

The Fischer-Clifford matrix $M(g)$ (see Figure 2) is partitioned row-wise into blocks $M_i(g)$, where each block corresponds to an inertia group \overline{H}_i . We write $|C_{\overline{G}}(x_j)|$, for each $x_j \in X(g)$, at the top of the columns of $M(g)$ and at the bottom we write $m_j = [C_g : C_{\overline{G}}(\overline{x}_j)] = |P| \frac{|C_G(g)|}{|C_{\overline{G}}(\overline{x}_j)|} = \frac{f_j |P|}{k}$. On the left of each row we write $|C_{H_i}(y_k)|$, where $[y_k]$, $k = 1, 2, \dots, r$, are the classes of an inertia factor H_i that fuse into the class $[g]$ of G . Since $|X(g)| = |R(g)|$ it follows that $M(g)$ is a square matrix of size $c(g)$. When there is no class fusion of an inertia factor H_i into a class $[g]$, the block $M_i(g)$ is omitted from $M(g)$.

$$M(g) = \begin{matrix} & |C_{\overline{G}}(x_1)| & |C_{\overline{G}}(x_2)| & \cdots & |C_{\overline{G}}(x_{c(g)})| \\ \begin{matrix} |C_G(g)| \\ |C_{H_2}(y_1)| \\ |C_{H_2}(y_2)| \\ \vdots \\ |C_{H_3}(y_1)| \\ |C_{H_3}(y_3)| \\ \vdots \\ |C_{H_4}(y_1)| \\ |C_{H_4}(y_2)| \\ \vdots \end{matrix} & \begin{pmatrix} 1 & 1 & \cdots & 1 \\ \hline \frac{|C_G(g)|}{|C_{H_2}(y_1)|} & a_{(2,y_1)}^2 & \cdots & a_{(2,y_1)}^{c(g)} \\ \frac{|C_G(g)|}{|C_{H_2}(y_2)|} & a_{(2,y_2)}^2 & \cdots & a_{(2,y_2)}^{c(g)} \\ \vdots & \vdots & \vdots & \vdots \\ \hline \frac{|C_G(g)|}{|C_{H_3}(y_1)|} & a_{(3,y_1)}^2 & \cdots & a_{(3,y_1)}^{c(g)} \\ \frac{|C_G(g)|}{|C_{H_3}(y_2)|} & a_{(3,y_2)}^2 & \cdots & a_{(3,y_2)}^{c(g)} \\ \vdots & \vdots & \vdots & \vdots \\ \hline \frac{|C_G(g)|}{|C_{H_4}(y_1)|} & a_{(4,y_1)}^2 & \cdots & a_{(4,y_1)}^{c(g)} \\ \frac{|C_G(g)|}{|C_{H_4}(y_2)|} & a_{(4,y_2)}^2 & \cdots & a_{(4,y_2)}^{c(g)} \\ \vdots & \vdots & \vdots & \vdots \end{pmatrix} \\ m_1 & m_2 & \cdots & m_{c(g)} \end{matrix}$$

Figure 2: The Fischer-Clifford Matrix $M(g)$

Instead of using the above formal definition, i.e. Equation 1, we use the arithmetical properties 1-8 of $M(g)$ below (see [12]) to compute the entries $a_{(i,y_k)}^j$. Since \overline{G} is a split extension of an elementary abelian group P of order 128, we use relation 6 to compute the values for column 1 of Figure 2. The relations 3 and 4 are inherited from the row and orthogonality relations of an ordinary character table of a finite group.

1. $a_{(1,g)}^j = 1$ for all $j = \{1, 2, \dots, c(g)\}$.
2. $|X(g)| = |R(g)|$.
3. $\sum_{j=1}^{c(g)} m_j a_{(i,y_k)}^j \overline{a_{(i',y'_k)}^j} = \delta_{(i,y_k),(i',y'_k)} \frac{|C_G(g)|}{|C_{H_i}(y_k)|} |P|$.
4. $\sum_{(i,y_k) \in R(g)} a_{(i,y_k)}^j \overline{a_{(i,y_k)}^{j'}} |C_{H_i}(y_k)| = \delta_{jj'} |C_{\overline{G}}(x_j)|$.
5. $M(g)$ is square and nonsingular.

Since P is an elementary abelian 2-group, then we obtain the additional properties 6-8 of $M(g)$ below,

6. $a_{(i,y_k)}^1 = \frac{|C_G(g)|}{|C_{H_i}(y_k)|}$.
7. $|a_{(i,y_k)}^1| \geq |a_{(i,y_k)}^j|$.
8. $a_{(i,y_k)}^j \equiv a_{(i,y_k)}^1 \pmod{2}$.

3.2. Computing the Fischer Clifford Matrices of \overline{G}

Using Table 2 and the construction process of a matrix $M(g)$ given in Section 3.1, we compute all the Fischer-Clifford matrices of \overline{G} and they are contained in Table 4.

Table 4: The Fischer-Clifford Matrices of \overline{G}

$M(g)$	$M(g)$
$M(1A) = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 28 & -28 & 4 & -4 \\ 36 & -36 & -4 & 4 \\ 63 & 63 & -1 & -1 \end{pmatrix}$	$M(2A) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 \\ 4 & -4 & 4 & -4 & 0 \\ 12 & -12 & -4 & 4 & 0 \\ 3 & 3 & 3 & 3 & -1 \\ 12 & 12 & -4 & -4 & 0 \end{pmatrix}$
$M(2B) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 & 1 \\ 4 & -4 & -4 & 4 & 0 & 0 \\ 4 & -4 & 4 & -4 & 0 & 0 \\ 1 & 1 & -1 & -1 & -1 & 1 \\ 3 & 3 & -3 & -3 & 1 & -1 \\ 3 & 3 & 3 & 3 & -1 & -1 \end{pmatrix}$	$M(3A) = \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$
$M(3B) = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & -1 & 1 & -1 \\ 3 & -3 & -1 & 1 \\ 3 & 3 & -1 & -1 \end{pmatrix}$	$M(4A) = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 4 & -4 & 0 & 0 \\ 1 & 1 & 1 & -1 \\ 2 & 2 & -2 & 0 \end{pmatrix}$
$M(4B) = \begin{pmatrix} 1 & 1 & 1 \\ 4 & -4 & 0 \\ 3 & 3 & -1 \end{pmatrix}$	$M(4C) = \begin{pmatrix} 1 & 1 & 1 \\ 4 & -4 & 0 \\ 3 & 3 & -1 \end{pmatrix}$
$M(6A) = \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$	$M(6B) = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & -1 & -1 & 1 \\ 1 & -1 & 1 & -1 \\ 1 & 1 & -1 & -1 \end{pmatrix}$
$M(7A) = \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$	$M(8A) = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & -1 & -1 \\ 1 & -1 & 1 & -1 \\ 1 & -1 & -1 & 1 \end{pmatrix}$
$M(8B) = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & -1 & -1 \\ 1 & -1 & 1 & -1 \\ 1 & -1 & -1 & 1 \end{pmatrix}$	$M(12A) = \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$
$M(12B) = \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$	$M(12C) = \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$

4. The Character Table of \overline{G}

Since we obtained the fusion maps of classes of the inertia factors H_i into G , the character tables of the inertia factors H_i (stored in the GAP library) and the Fischer-Clifford matrices of \overline{G} , we proceed to construct the ordinary character table of \overline{G} . A partial character table (see Figure 3) of \overline{G} is constructed by multiplying the columns $C_i(g)$ for $i \in \{1, 2, 3, 4\}$ of the ordinary character tables of the inertia factors H_i associated with the classes y_k of H_i , that are conjugate to $[g]_G$, by the rows of the Fischer-Clifford matrices in a block $M_i(g)$.

$$\begin{bmatrix} C_1(g) M_1(g) \\ C_2(g) M_2(g) \\ C_3(g) M_3(g) \\ C_4(g) M_4(g) \end{bmatrix}.$$

Figure 3: Partial Character Table of \overline{G}

To construct the set $\text{Irr}(G)$, we append the partial character tables coming from each class $[g_i]$, $i = 1, 2, \dots, 16$, of G as depicted in Figure 4.

$$\begin{bmatrix} C_1(g_1) M_1(g_1) & C_1(g_2) M_1(g_2) & \cdots & C_1(g_{16}) M_1(g_{16}) \\ C_2(g_1) M_2(g_1) & C_2(g_2) M_2(g_2) & \cdots & C_2(g_{16}) M_2(g_{16}) \\ C_3(g_1) M_3(g_1) & C_3(g_2) M_3(g_2) & \cdots & C_3(g_{16}) M_3(g_{16}) \\ C_4(g_1) M_4(g_1) & C_4(g_2) M_4(g_2) & \cdots & C_4(g_{16}) M_4(g_{16}) \end{bmatrix}.$$

Figure 4: Structure of $\text{Irr}(\overline{G})$

The character table of \overline{G} (see Table 5) obtained in this manner is a 53×53 -complex valued square matrix with the irreducible characters partitioned into blocks Δ_i , for $i \in \{1, 2, 3, 4\}$, such that $\Delta_1 = \{\chi_i | 1 \leq i \leq 16\}$, $\Delta_2 = \{\chi_i | 17 \leq i \leq 30\}$, $\Delta_3 = \{\chi_i | 31 \leq i \leq 39\}$ and $\Delta_4 = \{\chi_i | 40 \leq i \leq 53\}$, where $\chi_i \in \text{Irr}(\overline{G})$. The fusion map of the classes of \overline{G} into $Sp_8(2)$ is obtained with the aid of the GAP function ‘‘PossibleClassFusions’’ and the ordinary characters tables of \overline{G} and $Sp(8)$ and is captured in the last column of Table 3. A GAP routine in [19] was implemented for checking the consistency and accuracy of the character table of \overline{G} . To reconstruct the character table in GAP, that

is Table 5, interested readers can use the link below. The file contains the class orders, centralizer orders and the irreducible characters of \bar{G} which were computed in this paper.

https://drive.google.com/file/d/1pv7JhGzdiGJ1eB0h39WWEvmd_L0h5pte/view?usp=drive_link

Table 5: The Character Table of \bar{G}

$[g]_{\bar{G}}$	1A				2A					2B					3A		
$[x]_{\bar{G}}$	1A	2A	2B	2C	2D	2E	2F	2G	4A	2H	4B	2I	4C	4D	4E	3A	6A
χ_1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
χ_2	1	1	1	1	1	1	1	1	1	-1	-1	-1	-1	-1	-1	1	1
χ_3	6	6	6	6	-2	-2	-2	-2	-2	0	0	0	0	0	0	-3	-3
χ_4	6	6	6	6	-2	-2	-2	-2	-2	0	0	0	0	0	0	-3	-3
χ_5	7	7	7	7	-1	-1	-1	-1	-1	1	1	1	1	1	1	-2	-2
χ_6	7	7	7	7	-1	-1	-1	-1	-1	-1	-1	-1	-1	-1	-1	-2	-2
χ_7	14	14	14	14	-2	-2	-2	-2	-2	2	2	2	2	2	2	5	5
χ_8	14	14	14	14	-2	-2	-2	-2	-2	-2	-2	-2	-2	-2	-2	5	5
χ_9	14	14	14	14	6	6	6	6	6	0	0	0	0	0	0	-4	-4
χ_{10}	21	21	21	21	5	5	5	5	5	-3	-3	-3	-3	-3	-3	3	3
χ_{11}	21	21	21	21	5	5	5	5	5	3	3	3	3	3	3	3	3
χ_{12}	27	27	27	27	3	3	3	3	3	3	3	3	3	3	3	0	0
χ_{13}	27	27	27	27	3	3	3	3	3	-3	-3	-3	-3	-3	-3	0	0
χ_{14}	42	42	42	42	2	2	2	2	2	0	0	0	0	0	0	6	6
χ_{15}	56	56	56	56	-8	-8	-8	-8	-8	0	0	0	0	0	0	2	2
χ_{16}	64	64	64	64	0	0	0	0	0	0	0	0	0	0	0	-8	-8
χ_{17}	28	-28	-4	4	4	-4	4	-4	0	4	-4	-4	4	0	0	1	-1
χ_{18}	28	-28	-4	4	4	-4	4	-4	0	-4	4	4	-4	0	0	1	-1
χ_{19}	28	-28	-4	4	4	-4	4	-4	0	-4	4	4	-4	0	0	1	-1
χ_{20}	28	-28	-4	4	4	-4	4	-4	0	4	-4	-4	4	0	0	1	-1
χ_{21}	56	-56	-8	8	8	-8	8	-8	0	0	0	0	0	0	0	2	-2
χ_{22}	56	-56	-8	8	-8	8	-8	8	0	0	0	0	0	0	0	2	-2
χ_{23}	56	-56	-8	8	-8	8	-8	8	0	0	0	0	0	0	0	2	-2
χ_{24}	168	-168	-24	24	-8	8	-8	8	0	0	0	0	0	0	0	-3	3
χ_{25}	168	-168	-24	24	-8	8	-8	8	0	0	0	0	0	0	0	-3	3
χ_{26}	168	-168	-24	24	-8	8	-8	8	0	0	0	0	0	0	0	-3	3
χ_{27}	168	-168	-24	24	-8	8	-8	8	0	0	0	0	0	0	0	-3	3
χ_{28}	224	-224	-32	32	0	0	0	0	0	-8	8	8	-8	0	0	8	-8
χ_{29}	224	-224	-32	32	0	0	0	0	0	8	-8	-8	8	0	0	8	-8
χ_{30}	336	-336	-48	48	16	-16	16	-16	0	0	0	0	0	0	0	-6	6
χ_{31}	36	-36	4	-4	12	-12	-4	4	0	4	4	-4	-4	0	0	0	0
χ_{32}	36	-36	4	-4	12	-12	-4	4	0	-4	-4	4	4	0	0	0	0
χ_{33}	216	-216	24	-24	-24	24	8	-8	0	0	0	0	0	0	0	0	0
χ_{34}	216	-216	24	-24	24	-24	-8	8	0	0	0	0	0	0	0	0	0
χ_{35}	216	-216	24	-24	24	-24	-8	8	0	0	0	0	0	0	0	0	0
χ_{36}	252	-252	28	-28	-12	12	4	-4	0	-4	-4	4	4	0	0	0	0
χ_{37}	252	-252	28	-28	-12	12	4	-4	0	4	4	-4	-4	0	0	0	0
χ_{38}	288	-288	32	-32	0	0	0	0	0	-8	-8	8	8	0	0	0	0
χ_{39}	288	-288	32	-32	0	0	0	0	0	8	8	-8	-8	0	0	0	0
χ_{40}	63	63	-1	-1	15	15	-1	-1	-1	7	-1	7	-1	-1	-1	0	0
χ_{41}	63	63	-1	-1	-9	-9	7	7	-1	1	5	1	5	-3	1	0	0
χ_{42}	63	63	-1	-1	-9	-9	7	7	-1	-1	-5	-1	-5	3	-1	0	0
χ_{43}	63	63	-1	-1	15	15	-1	-1	-1	-7	1	-7	1	1	1	0	0
χ_{44}	126	126	-2	-2	6	6	6	6	-2	8	4	8	4	-4	0	0	0
χ_{45}	126	126	-2	-2	6	6	6	6	-2	-8	-4	-8	-4	4	0	0	0
χ_{46}	189	189	-3	-3	-3	-3	13	13	-3	-3	-3	-3	-3	-3	5	0	0
χ_{47}	189	189	-3	-3	-3	-3	13	13	-3	3	3	3	3	3	-5	0	0
χ_{48}	189	189	-3	-3	21	21	5	5	-3	-3	9	-3	9	1	-3	0	0
χ_{49}	189	189	-3	-3	21	21	5	5	-3	3	-9	3	-9	-1	3	0	0
χ_{50}	378	378	-6	-6	-30	-30	2	2	2	0	0	0	0	0	0	0	0
χ_{51}	378	378	-6	-6	-6	-6	-6	-6	2	-6	6	-6	6	-2	2	0	0
χ_{52}	378	378	-6	-6	-6	-6	-6	-6	2	6	-6	6	-6	2	-2	0	0
χ_{53}	378	378	-6	-6	18	18	-14	-14	2	0	0	0	0	0	0	0	0

$[g]_G$	3B				4A				4B			4C			6A	
$[x]_{\overline{G}}$	3B	6B	6c	6D	4F	4G	4H	4I	4J	4K	4L	4M	4N	4O	6E	6F
χ_1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
χ_2	1	1	1	1	1	1	1	-1	-1	-1	1	1	1	1	1	1
χ_3	0	0	0	0	-2	-2	-2	0	0	0	2	2	2	2	1	1
χ_4	0	0	0	0	-2	-2	-2	0	0	0	2	2	2	2	1	1
χ_5	1	1	1	1	3	3	3	-3	-3	-3	-1	-1	-1	-1	2	2
χ_6	1	1	1	1	3	3	3	3	3	3	-1	-1	-1	-1	2	2
χ_7	-1	-1	-1	-1	2	2	2	-2	-2	-2	2	2	2	2	1	1
χ_8	-1	-1	-1	-1	2	2	2	2	2	2	2	2	2	2	1	1
χ_9	2	2	2	2	-2	-2	-2	0	0	0	2	2	2	2	0	0
χ_{10}	0	0	0	0	1	1	1	1	1	1	1	1	1	1	-1	-1
χ_{11}	0	0	0	0	1	1	1	-1	-1	-1	1	1	1	1	-1	-1
χ_{12}	0	0	0	0	3	3	3	3	3	3	-1	-1	-1	-1	0	0
χ_{13}	0	0	0	0	3	3	3	-3	-3	-3	-1	-1	-1	-1	0	0
χ_{14}	0	0	0	0	-6	-6	-6	0	0	0	-2	-2	-2	-2	2	2
χ_{15}	2	2	2	2	0	0	0	0	0	0	0	0	0	0	-2	-2
χ_{16}	-2	-2	-2	-2	0	0	0	0	0	0	0	0	0	0	0	0
χ_{17}	1	-1	1	-1	4	-4	0	4	-4	0	0	0	0	0	1	-1
χ_{18}	1	-1	1	-1	4	-4	0	-4	4	0	0	0	0	0	1	-1
χ_{19}	1	-1	1	-1	4	-4	0	4	-4	0	0	0	0	0	1	-1
χ_{20}	1	-1	1	-1	4	-4	0	-4	4	0	0	0	0	0	1	-1
χ_{21}	2	-2	2	-2	-8	8	0	0	0	0	0	0	0	0	2	-2
χ_{22}	2	-2	2	-2	0	0	0	0	0	0	0	0	0	0	-2	2
χ_{23}	2	-2	2	-2	0	0	0	0	0	0	0	0	0	0	-2	2
χ_{24}	0	0	0	0	8	-8	0	8	-8	0	0	0	0	0	1	-1
χ_{25}	0	0	0	0	8	-8	0	-8	8	0	0	0	0	0	1	-1
χ_{26}	0	0	0	0	-8	8	0	0	0	0	0	0	0	0	1	-1
χ_{27}	0	0	0	0	-8	8	0	0	0	0	0	0	0	0	1	-1
χ_{28}	-1	1	-1	1	0	0	0	0	0	0	0	0	0	0	0	0
χ_{29}	-1	1	-1	1	0	0	0	0	0	0	0	0	0	0	0	0
χ_{30}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	-2	2
χ_{31}	3	-3	-1	1	0	0	0	0	0	0	4	-4	0	0	0	0
χ_{32}	3	-3	-1	1	0	0	0	0	0	0	4	-4	0	0	0	0
χ_{33}	0	0	0	0	0	0	0	0	0	0	8	-8	0	0	0	0
χ_{34}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{35}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{36}	3	-3	-1	1	0	0	0	0	0	0	-4	4	0	0	0	0
χ_{37}	3	-3	-1	1	0	0	0	0	0	0	-4	4	0	0	0	0
χ_{38}	-3	3	1	-1	0	0	0	0	0	0	0	0	0	0	0	0
χ_{39}	-3	3	1	-1	0	0	0	0	0	0	0	0	0	0	0	0
χ_{40}	3	3	-1	-1	3	3	-1	3	3	-1	3	3	-1	-1	0	0
χ_{41}	3	3	-1	-1	3	3	-1	-3	-3	1	-1	-1	3	-1	0	0
χ_{42}	3	3	-1	-1	3	3	-1	3	3	-1	-1	-1	3	-1	0	0
χ_{43}	3	3	-1	-1	3	3	-1	-3	-3	1	3	3	-1	-1	0	0
χ_{44}	-3	-3	1	1	6	6	-2	0	0	0	2	2	2	-2	0	0
χ_{45}	-3	-3	1	1	6	6	-2	0	0	0	2	2	2	-2	0	0
χ_{46}	0	0	0	0	-3	-3	1	-3	-3	1	-3	-3	1	1	0	0
χ_{47}	0	0	0	0	-3	-3	1	3	3	-1	-3	-3	1	1	0	0
χ_{48}	0	0	0	0	-3	-3	1	-3	-3	1	1	1	-3	1	0	0
χ_{49}	0	0	0	0	-3	-3	1	3	3	-1	1	1	-3	1	0	0
χ_{50}	0	0	0	0	-6	-6	2	0	0	0	6	6	-2	-2	0	0
χ_{51}	0	0	0	0	6	6	-2	6	6	-2	-2	-2	-2	2	0	0
χ_{52}	0	0	0	0	6	6	-2	-6	-6	2	-2	-2	-2	2	0	0
χ_{53}	0	0	0	0	-6	-6	2	0	0	0	-2	-2	6	-2	0	0

$[g]_G$	6B				7A		8A				8B				12A		12B		12C	
$[x]_{\bar{G}}$	6G	6H	12A	12B	7A	14A	8A	8B	8C	8D	8E	8F	8G	8H	12C	12D	12E	12F	12G	12h
χ_1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
χ_2	-1	-1	-1	-1	1	1	-1	-1	-1	-1	1	1	1	1	1	1	-1	-1	-1	-1
χ_3	0	0	0	0	-1	-1	0	0	0	0	0	0	0	0	1	1	C	C	-C	-C
χ_4	0	0	0	0	-1	-1	0	0	0	0	0	0	0	0	1	1	-C	-C	C	C
χ_5	1	1	1	1	0	0	-1	-1	-1	-1	-1	-1	-1	-1	0	0	0	0	0	0
χ_6	-1	-1	-1	-1	0	0	1	1	1	1	-1	-1	-1	-1	0	0	0	0	0	0
χ_7	-1	-1	-1	-1	0	0	0	0	0	0	0	0	0	0	-1	-1	1	1	1	1
χ_8	1	1	1	1	0	0	0	0	0	0	0	0	0	0	-1	-1	-1	-1	-1	-1
χ_9	0	0	0	0	0	0	0	0	0	0	0	0	0	0	-2	-2	0	0	0	0
χ_{10}	0	0	0	0	0	0	-1	-1	-1	-1	-1	-1	-1	-1	1	1	1	1	1	1
χ_{11}	0	0	0	0	0	0	1	1	1	1	-1	-1	-1	-1	1	1	-1	-1	-1	-1
χ_{12}	0	0	0	0	-1	-1	-1	-1	-1	-1	1	1	1	1	0	0	0	0	0	0
χ_{13}	0	0	0	0	-1	-1	1	1	1	1	1	1	1	1	0	0	0	0	0	0
χ_{14}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{15}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{16}	0	0	0	0	1	1	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{17}	1	-1	-1	1	0	0	0	0	0	0	2	-2	0	0	1	-1	1	-1	1	-1
χ_{18}	-1	1	1	-1	0	0	0	0	0	0	2	-2	0	0	1	-1	-1	1	-1	1
χ_{19}	-1	1	1	-1	0	0	0	0	0	0	-2	2	0	0	1	-1	1	-1	1	-1
χ_{20}	1	-1	-1	1	0	0	0	0	0	0	-2	2	0	0	1	-1	-1	1	-1	1
χ_{21}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	-2	2	0	0	0	0
χ_{22}	0	0	0	0	0	0	0	0	0	0	0	0	B	-B	0	0	0	0	0	0
χ_{23}	0	0	0	0	0	0	0	0	0	0	0	0	-B	B	0	0	0	0	0	0
χ_{24}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	-1	1	-1	1	-1	1
χ_{25}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	-1	1	1	-1	1	-1
χ_{26}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	1	-1	C	-C	-C	C
χ_{27}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	1	-1	-C	C	C	-C
χ_{28}	1	-1	-1	1	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{29}	-1	1	1	-1	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{30}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{31}	1	-1	1	-1	1	-1	2	-2	0	0	0	0	0	0	0	0	0	0	0	0
χ_{32}	-1	1	-1	1	1	-1	-2	2	0	0	0	0	0	0	0	0	0	0	0	0
χ_{33}	0	0	0	0	-1	1	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{34}	0	0	0	0	-1	1	0	0	A	-A	0	0	0	0	0	0	0	0	0	0
χ_{35}	0	0	0	0	-1	1	0	0	-A	A	0	0	0	0	0	0	0	0	0	0
χ_{36}	-1	1	-1	1	0	0	2	-2	0	0	0	0	0	0	0	0	0	0	0	0
χ_{37}	1	-1	1	-1	0	0	-2	2	0	0	0	0	0	0	0	0	0	0	0	0
χ_{38}	1	-1	1	-1	1	-1	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{39}	-1	1	-1	1	1	-1	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{40}	1	1	-1	-1	0	0	1	1	-1	-1	1	1	-1	-1	0	0	0	0	0	0
χ_{41}	1	1	-1	-1	0	0	-1	-1	1	1	-1	-1	1	1	0	0	0	0	0	0
χ_{42}	-1	-1	1	1	0	0	1	1	-1	-1	-1	-1	1	1	0	0	0	0	0	0
χ_{43}	-1	-1	1	1	0	0	-1	-1	1	1	1	1	-1	-1	0	0	0	0	0	0
χ_{44}	-1	-1	1	1	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{45}	1	1	-1	-1	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{46}	0	0	0	0	0	0	1	1	-1	-1	1	1	-1	-1	0	0	0	0	0	0
χ_{47}	0	0	0	0	0	0	-1	-1	1	1	1	1	-1	-1	0	0	0	0	0	0
χ_{48}	0	0	0	0	0	0	1	1	-1	-1	-1	-1	1	1	0	0	0	0	0	0
χ_{49}	0	0	0	0	0	0	-1	-1	1	1	-1	-1	1	1	0	0	0	0	0	0
χ_{50}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{51}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{52}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
χ_{53}	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0

$A = -2\sqrt{2}i, B = -2\sqrt{2}, C = -\sqrt{3}i$

Data availability statement

The authors confirm that the data supporting the findings of this study are available within the article.

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